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REMARKS ON NUCLEAR FISSION N. J. Carrel 2-1-5-6/45TM By REPORT LIBRARY M. Markin

General character of fission yield curves derived from formula:

Characteristic dependance of  $\lceil R \rceil$ ,  $\lceil R \rceil$  and  $\lceil R \rceil$  indicated in Figure 1. Assuming a simple course of curve for 50 as indicated in Figure 2 one obtains curves for  $\sigma_{\rm p}$  of the types illustrated on same figure.

Curve or (1) would correspond to neutron fission of 49. Whether peculiar dip in  $\mathcal{T}_{p}^{49}$  is primarily due to competition between neutron escape and fission or whether it may be partly due to some marked minimum of ce in this region can of course not be decided apriori; but it is interesting that a curve for  $\sigma_{\alpha}$  of the type assumed in a qualitative way explains not only the experimental curve for  $\sigma_{\rm r}^{49}$  but also for  $\sigma_{\rm r}^{25}$ . In the latter case the observation that  $\sigma_{\rm F}$  in the region (0.1 Mev < E<sub>W</sub> < 0.5 Mev) is closely proportional to 1/v might in fact be the combined result of the gradual decrease of the  $\sigma_{\sigma}$  curve and the competition between  $\Gamma_{F}$  and  $\Gamma_{N}$  which in this region

is rapidly increasing. The observed flat part of the  $\sigma_{\rm F}^{25}$  curve for high energies should thus correspond to the later more parallel course of the

curves for F and I.

While  $\mathcal{T}_{F}^{(3)}$  qualitatively corresponds to  $\mathcal{T}_{F}^{28}$  it is obviously

necessary to discuss the curves for  $\Gamma_{\rm F}$  and  $\Gamma_{\rm N}$  in somewhat more detail in order to explain the observed peculiar differences between  $\sigma_{
m F}^{28}$  and  $\sigma_{
m F}^{02}$  indi-

cated by the dotted lines in Figure 3. It is here assumed that the  $\Gamma_{\rm N}$ 

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curve has some characteristic irregularities due to the successive appearance of levels for inelastic scattering. The  $\Gamma_F^{28}$  and  $\Gamma_F^{02}$  are assumed apart from a displacement of 0.2 MeV to be essentially similar since the short lifetime of the transition state will have the effect almost entirely to efface any specific level structure.

2. Estimate of critical fission energy F and neutron binding energy W

On the assumption that

$$F = f(Z^2/A)$$

we may in the comparatively small region of Z2/A concerned write

$$F = a - b(Z^2/A)$$

A new possibility of estimating a and b is given by the measurements of the threshold for photo fission. Taking F = 5.7 MeV for 28 and F = 6.1 MeV for 02 one gets

$$F = -0.52(Z - 92) + 0.1(A - 238)$$

The values derived from this interpolation formula are given in the second column of the Table on page 3 where the numbers in the first column refer to the values of Z and A of the compound system. In the third column are given the values for F - W which in the case of the 63 first isotopes are roughly estimated by comparison with the curves on Figure 1 and 2. The resultant values of W given in the fourth column would seem quite reasonable. In particular the difference in W for nuclei of even and odd values of A - Z is of the order of magnitude to be expected from general evidence as regards the radioactive properties of such nuclei.



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z <sup>A</sup>	F	F - W	W
<sub>Մ</sub> 234	5 <b>•3</b>	≈ - 1	6.3
Pu <sup>240</sup>	4.9	<b>≈ -</b> 1	5.9
<sub>U</sub> 236	5.5	≈ - ½	6.0
<sub>Մ</sub> 235	5 <b>.4</b>	≈ + 0.3	5.1
Np238	5.2	+ 0.4	4.8
Pa 232	5.6	. 0.5	5.1
<sub>U</sub> 239	5.8	1.0	4.8
Th <sup>233</sup>	6.2	1.2	5.0

## 3. Comparison between level spectra of 28 and 25 at low energies.

The discussion is based on the disperson formulae

$$\sigma_{\rm R} = \pi \, \tilde{\chi}^2 \cdot \frac{\Gamma_{\rm N} \, \Gamma_{\rm R}}{(E - E_{\rm O})^2 + (\Gamma/2)^2} \; ; \; \sigma_{\rm F} = \pi \, \tilde{\chi}^2 \cdot \frac{\Gamma_{\rm N} \, \Gamma_{\rm F}}{(E - E_{\rm O})^2 + (\Gamma/2)^2}$$

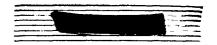
Since in 28 we have  $\Gamma_F = 0$  and  $\Gamma_N < \Gamma_R$ , the width of the lines will be determined by  $\Gamma_R$  and since in 25,  $\Gamma_F > \Gamma_R$  the width will in this case be determined by  $\Gamma_F$ . For both cases we have, therefore,

$$B = \int ddE = 4\pi \lambda^2 \Gamma_N \sim 1/v$$

From the curvo for  $\sigma_F^{25}$  in Fig. 4 one finds that Bv for the levels below 10 ev is almost 100 times smaller than for the wellknown level at 7 ev in 28. Assuming that, as indicated in Fig. 4, the spectrum in 25 consists of both weak and strong lines with a different level spacing, the latter levels are found to have values of Bv comparable with the 28 line. UNCLASSIFIE

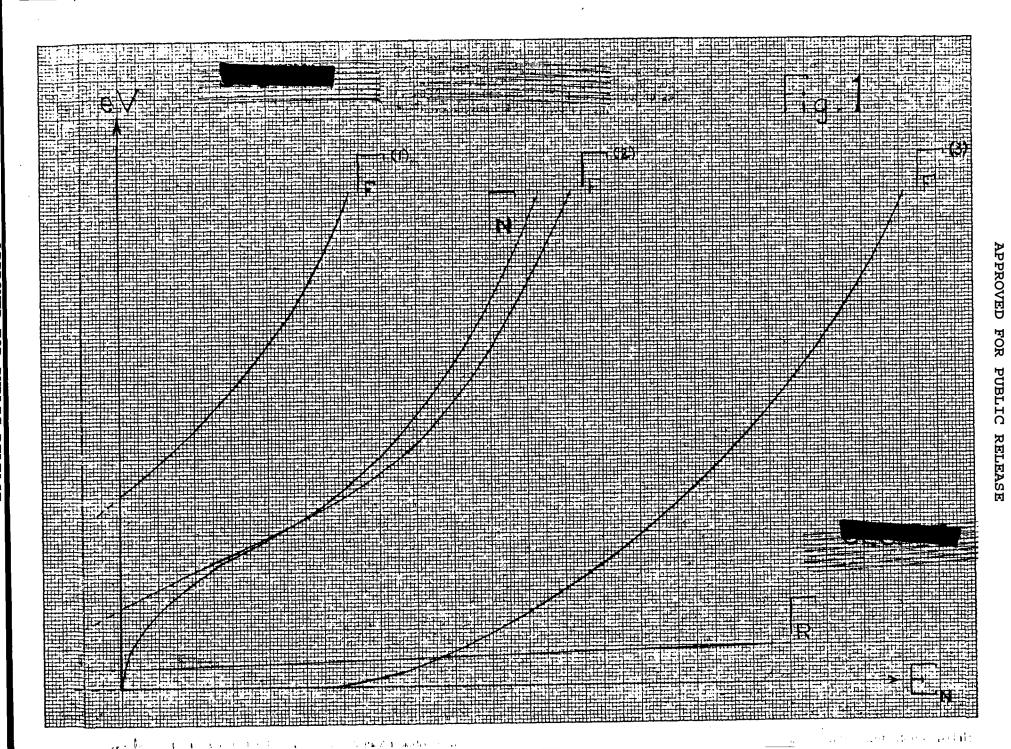
A suggested explanation of the appearance of the two level systems

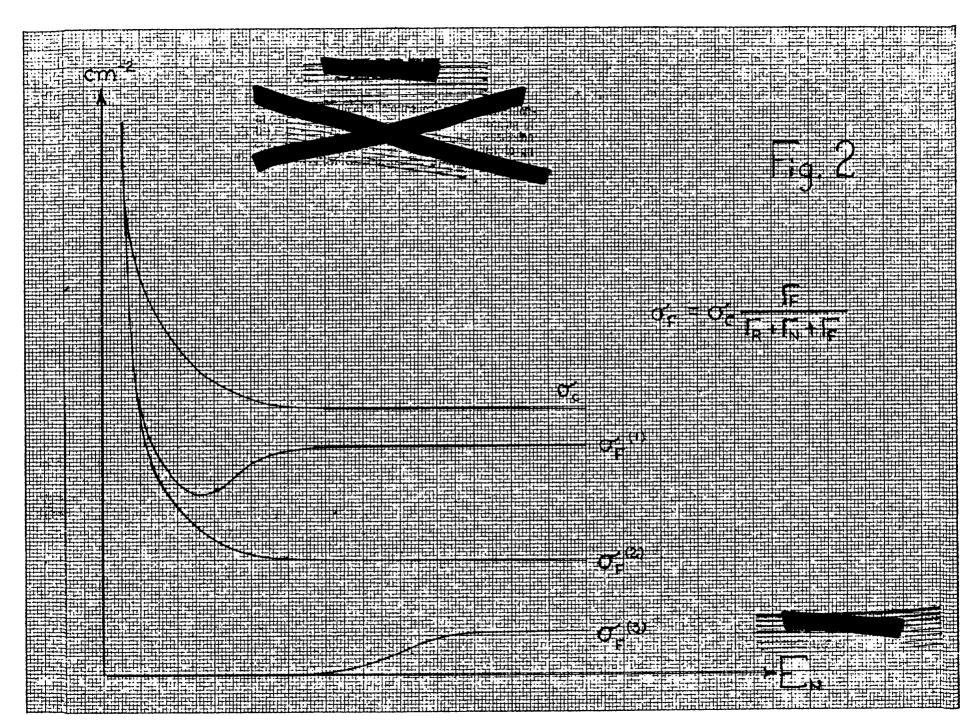
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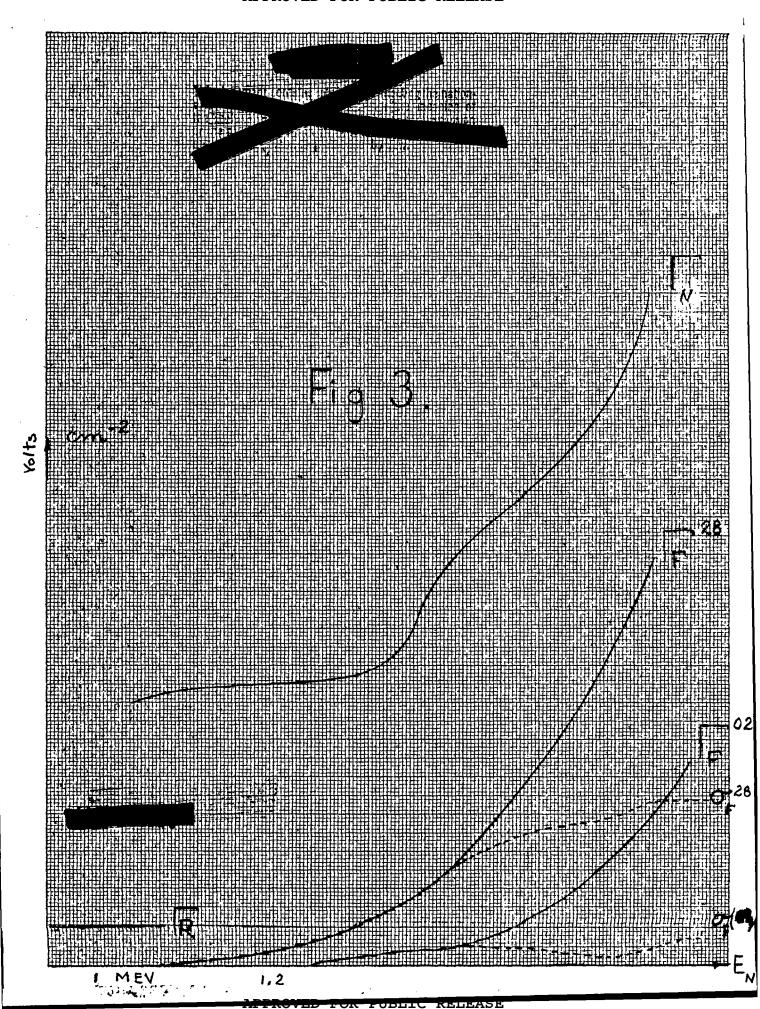


is schematically represented in Fig. 5 in which is also indicated a possible explanation of the remarkable increase in  $v \vec{\sigma_F}^{25}$  in the region below 100,000 eV, where the neutrons may be able to give angular momentum to the nucleus with the result that the difficulty in forming levels of the compound state with spin S-1 may disappear. This suggestion has given rise to discussions about impact of neutrons with  $l \neq 0$  which have not yet been concluded and which are hoped to show whether the suggested peculiar increase in  $\Gamma_{N\overline{V}}^{-1}$  within the S-1 level system actually does take place.

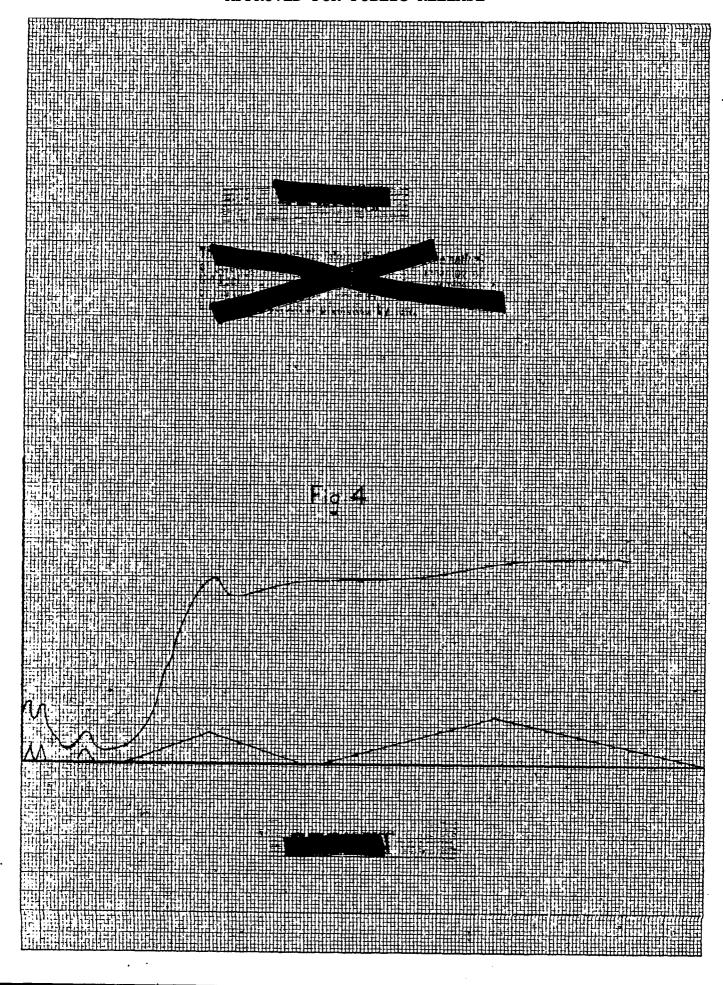








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